

Optimization and Limitations of a Multi-Level Adaptive-Order Fast Multipole Algorithm for Acoustical Calculations

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Abstract: The Multi-Level Fast Multipole Method (MLFMM) allows the computation of acoustical problems based on the Boundary Element Method (BEM) where the discretized models of the corresponding structures may consist of a huge number of elements. The required calculation time and the memory requirements are much less when compared with conventional methods because the algorithm uses a level-based composition of the potentials from different point sources to acoustic multipoles, which highly accelerates the computation of the matrix-vector-products required for iterative solvers. A multi-level single-order variation of the algorithm developed during a previous research project was extended to a multi-level adaptive-order version, which was analyzed and optimized with respect to quality, performance and parallelization issues. The insights gained will be presented using different test cases and the results achieved will be compared with analytical solutions and results based on conventional BEM- and FEM-based calculations. **Keywords:** Fast Multipole Method, Optimization, Numerical Methods

1. Introduction

The Multi-Level Fast Multipole Method (MLFMM) describes a fast algorithm to accelerate the matrix-vector product which is required for the iterative solution of BEM-based calculations without ever assembling the complete matrix. The method is suited for big problems where the interactions between huge numbers of source (N_x) and destination points (N_y) must be considered.



Figure 1 - direct interactions (conventional BEM)

Figure 2 - cluster based interactions (MLFMM)

The decrease of the number of interactions when using the multi-level version of the algorithm can be identified by comparing Figure 1 and Figure 2.

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The effects of all individual sources x_i within a source cluster C_x of radius R are summarized in a multipole source at z_x . Their potential is transformed to a remote target cluster C_y centered at z_y and distributed to the target points y_j (Figure 3).



Figure 3 - Decomposition of the path between source and destination points of two clusters (based on [Sak4], p. 347, Fig. 12.9)

For each part of the path between source and destination points a corresponding transfer function $(h, \mu^M \text{ and } f)$ is used as shown in Figure 4, the full details and formulations may be found in [Sak4].



Figure 4 - path dependent transfer functions h, μ^M and f

One if the most critical functions here is the translation operator μ^M which "transfers" the multipole potential from a source cluster (the far field signature F_s) to the center of a target cluster with the near field signature N^s . This operator can be represented as a truncated series (Eq. (1)) with the maximum order set to O_{mp} ("multipole order"):

$$\mu^{M}(\vec{a},\hat{s}) = \frac{1}{4\pi} \sum_{l=1}^{O_{mp}} (2l+1) \,\mathbf{i}^{l} \,h_{l}(k \,|\vec{a}|) \,P_{l}(\hat{s} \cdot \hat{a}) \tag{1}$$

with \vec{a} distance vector between cluster centers

- \hat{s} set of vectors on the unit sphere
- *k* wave number

 $h_l(k |\vec{a}|)$ Hankel function

 $P_l(\hat{s}\cdot\hat{a})$ Legendre polynomials

There are many articles in literature dealing with the determination of the parameters used [Sak6, Yas7] and the specialties of the Hankel function h_l and Legendre polynoms P_l , but this would exceed the framework of this paper.

First we implemented a fixed-order version (Vs. 1.x) of the MLFMM algorithm [Bur1, Bur2] based on [Sak4] to check the limits with regard to element numbers and frequency. Then the algorithm was extended to an adaptive-level version (Vs. 2.x) using [Che5, Sak6, Che8] where the multipole order is defined as a function $O_{mp}(k, |\vec{a}|)$ and calculated at runtime. In addition, other optimizations to reduce the solution time were made, especially within the precalculation of geometry-based values.

2. Examples and results

2.1 Rigid sphere

To test the hardware dependent limits of the implemented MLFMM algorithm, a simple rigid sphere (r = 0.5 m, Figure 5) placed in water and hit by a plane wave using a frequency of 1 kHz was discretized with resulting element numbers between 1.000 and 5 millions.



Figure 5 - rigid sphere

The resulting solving times of the fixed-order MLFMM algorithm (Vs. 1.020 and 1.030) are shown in Table 1 and Figure 6, using a fixed multipole order of $O_{mp} = 6$ and an iteration error of $e_{iter} \le 10^{-10}$. All calculations ran on a 12 core workstation with 48 GB RAM, thus limiting the number of unknowns to $\approx 64,000$ for a full matrix.

number of elements	direct solver	iterative solver (GMRES)		iterative solver w/MLFMM, Vs. 1.030			
N_{elem}	Δt_{solve}	Δt_{solve}	N_{iter}	Δt_{solve}	N_{iter}	Δt_{MVP1}	$\oslash \Delta t_{MVP}$
1k	0.20 s	0.16 s	12	0.64 s	13	0.093 s	0.045 s
2.5k	1.23 s	0.73 s	11	0.62 s	12	0.141 s	0.042 s
5k	7.44 s	3.03 s	11	2.06 s	13	0.328 s	0.142 s
10k	43.99 s	11.92 s	11	2.59 s	12	0.562 s	0.180 s
20k	310.75 s	52.56 s	11	7.63 s	13	1.202 s	0.530 s
50k	4,352.43 s	329.37 s	11	15.13 s	12	3.433 s	1.053 s
100k	n.a.	11,710.95 [*] s	11	31.87 s	12	5.975 s	2.331 s
200k	n.a.	46,635.42 [*] s	11	57.16 s	12	12.823 s	3.989 s
500k	n.a.	n.a.		138.26 s	12	28.205 s	9.869 s
1M	n.a.	n.a.		373.61 s	12	68.500 s	27.596 s
2M	n.a.	n.a.		561.82 s	12	114.973 s	40.332 s
5M	n.a.	n.a.		1,922.26 s	13	315.325 s	133.179 s

Table 1 - solving times for the rigid sphere using different methods

(a full matrix was rebuild for each matrix-vector product for results marked with an ^{*})



Figure 6 - solving times for the rigid sphere using different methods

These solution times were significantly reduced using the adaptive-order version (Table 2, Figure 7).

elements	iterativ	ve solver w/M	multipole order			
N_{elem}	Δt_{solve}	N_{iter}	Δt_{MVP1}	$\oslash \Delta t_{MVP}$	$O_{mp,min}$	$O_{mp,max}$
1k	0.29	13	0.071	0.015	6	10
2.5k	0.42	13	0.120	0.024	6	10
5k	0.91	12	0.319	0.047	6	10
10k	1.33	12	0.560	0.067	6	10
20k	2.57	12	1.115	0.127	6	10
50k	7.38	12	4.037	0.291	6	10
100k	12.79	12	6.557	0.543	6	10
200k	25.88	12	13.630	1.083	6	10
500k	63.74	12	33.053	2.661	6	10
1M	134.61	12	71.360	5.491	6	10
2M	257.12	12	140.226	10.100	6	10
5M	856.73	12	532.744	25.987	6	10

Table 2 - solving times for the rigid sphere using different MLFMM versions



Figure 7 - solving times for the rigid sphere using different MLFMM versions

2.2 Ellipsoid

Another problem with the use of the fixed-order algorithm is that the solution due to the error in the matrix-vector products does not match the expected solution, although the iterative error was below a given value. To illustrate this problem, a rigid ellipsoid ($2 \times 4 \times 1$ m, $N_{elem} = 17,300$) placed in water is hit by a plane wave with at an incident angle of 30° using a frequency of f = 2.5 kHz.

Figure 8 shows the expected result using a conventional matrix-based BEM calculation.



Figure 8 - Ellipsoid, $abs(p_{surf})$, matrix-based

This model is difficult for the MLFMM due to the different dimensions which require high multipole orders for correct results due to larger cluster distances. Figure 9 and Figure 10 show the results when using the fixed-order MLFMM algorithm.



Figure 10 - Ellipsoid, $abs(p_{surf})$, MLFMM, fixed order of $O_{mp} = 15$

The dependence of the quality on the multipole order O_{mp} is clear. The adaptive-level version gives the best performance and quality (Figure 11).



Figure 11 - Ellipsoid, $abs(p_{surf})$, MLFMM, adaptive order $O_{mp} = 6 \dots 55$

3. Conclusions and future work

The results presented show that the adaptive and performance-optimized version of the MFLMM algorithm gives better quality and faster results.

At higher frequencies, the MLFMM shows significant qualitative differences, since the error in the method-based matrix-vector product forms a stronger effect. Further investigations to optimize the code for higher multipole expansion orders are necessary.

Likewise, a suitable preconditioner should be used to achieve a better convergence of the iterative method. First results were published in [Och3], but these were not really successful, so further work is needed here.

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More information can be found at the website of our research group "computational acoustics": http://projekt.beuth-hochschule.de/ca